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**Thermal Physics. Charles Kittel. Herbert**

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SOLUTIONS K&K 2.1 Part a Suppose  $g(U) =$   
 $C U^{3N-2}$ , where  $C$  is a constant and  $N$  is the  
number of particles. Show that  $U = \frac{3}{2} N \langle \epsilon \rangle$ . We  
use the definition of temperature as  $\frac{1}{\beta} = \left( \frac{\partial \ln g}{\partial U} \right)^{-1}$   
 $\frac{1}{\beta} = \left( \frac{\partial \ln g}{\partial U} \right)^{-1}$  So, let's calculate  $\beta$ .  $\beta = \ln(g) = 3N$   
 $\frac{1}{2} \ln(U) + \ln(C)$  (2) Therefore,  $\frac{1}{\beta} = \frac{3N}{2U}$   
 $= \frac{1}{\beta}$  (3) So finally,  $U = \frac{3}{2} N \langle \epsilon \rangle$  (4) Part b  
Show that  $\frac{1}{\beta} = \frac{1}{\beta}$

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The energy is  $E = 2 \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2 k^2}{m}$ .  
(b) At  $k = \frac{P}{\hbar}$  the determinantal equation is  $(\frac{P}{\hbar a}) \sin Ka + \cos Ka = -1$ . In the same limit this equation has solutions  $Ka = \frac{P}{\hbar a} + \frac{1}{2} \frac{P}{\hbar a}$ , where  $\frac{1}{2} \frac{P}{\hbar a} \ll 1$ . We expand to obtain  $(\frac{P}{\hbar a} + \frac{1}{2} \frac{P}{\hbar a}) \sin Ka + \cos Ka = -1$ , which has the solution  $\frac{1}{2} \frac{P}{\hbar a} = 0$  and  $\frac{1}{2} \frac{P}{\hbar a} = 2P/\hbar a$ .

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offers a modern approach to thermal physics that is based on the idea that all physical systems can be described in terms of their discrete quantum states, rather than drawing on 19th-century classical mechanics concepts

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Kittel and Kroemer teach thermal physics through the concepts of statistical mechanics and lay a strong foundation in this way of viewing physics. From there they introduce quantities such as the partition function, the free energy, Gibbs energy, entropy and so on. They do not take the engineering approach to thermal physics, such as btu's etc.

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